## **Short Communications**

## Force Field of Ethylthiocyanate\* ALF BJØRSETH\*\*

Department of Chemistry and Center for Structural Studies, The University of Texas at Austin, Austin, Texas 78712, U.S.A.

The vibrational spectra and conformation of ethylthiocyanate have been of interest to several research groups during the last years. A thorough analysis of the vibrational spectrum was done by Hirschman et al.,1 who interpreted the spectrum on the basis of two rotational isomers, anti (with  $C_s$  symmetry) and gauche ( $C_1$  symmetry). The anti-form was found to be the more stable. Crowder 2 calculated a force field which reproduces most of the observed frequencies, based on the anti-conformation. From a recent microwave investigation<sup>3</sup> of ethylthiocyanate, however, it was concluded that the gauche form was the more stable, and no lines originating from the anti conformer were assigned. This led to a reinvestigation 4 of the IR-spectrum in the solid (crystalline and amorphous), liquid, and vapour phase and the liquid Raman spectrum. This work concludes that all bands in all states can be ascribed to one isomer only, and that the polarization data and vapor phase contours indicate that this rotamer is the gauche isomer. In the present communication the determination of the force field based on the gauche conformation is reported.

Method of calculation. The normal coordinate program used in this work is described by Gwinn.<sup>5</sup> In this program mass weighted cartesian coordinates are used and rotation and translation are not separated until the matrix diagonalization is performed. The V matrix is generated numerically and only one transformation, from cartesian to normal coordinates,

is involved.

In the initial calculations a valence force field

$$\mathbf{V} = \frac{1}{2} \sum K_{ij} (\Delta r_{ij})^2 + \frac{1}{2} \sum H_{ijk} (\Delta \phi_{ijk})^2 + \frac{1}{2} \sum Y_{ijkl} (\Delta t_{ijkl})^2$$

was used, where the diagonal force constants K, H, and Y represent stretching, bending, and torsion, respectively. In many cases, redundant internal coordinates are used, i.e. the three  $H_{CCH}$  and three  $H_{HCH}$  of the methyl group. During the transformation to the basis of cartesian displacement coordinates these redundancies are removed automatically. Likewise, the torsional force constant is read for all HCCH combinations in the ethyl group, i.e.  $9Y_{
m HCCH}$  force constants (with the same value). To improve the agreement between observed and calculated frequencies it was found necessary to introduce interaction constants, defined as  $F^1$  for interaction between two stretching force constants,  $F^2$  between stretching and bending, and  $F^3$  between two bending force constants. The structure of the molecule was taken from the microwave work.3

Table 1. Force constants of CH<sub>3</sub>CH<sub>2</sub>SCN.

$\mathrm{Value}^a$	
16.73	
4.70	
0.392	
0.369	
0.68	
1.16	
0.66	
0.56	
0.66	
0.50	
0.54	
0.02	
0.10	
0.30	
0.51	
0.35	
-0.125	
0.015	
-0.015	
	0.392 0.369 0.68 1.16 0.66 0.56 0.50 0.54 0.02 0.0114 0.10 0.30 0.51 0.35 - 0.125 0.015

<sup>&</sup>lt;sup>a</sup>Units are mdyn/Å for K and  $F^1$ , mdyn for  $F^2$  and mdyn Å for H, Y,  $F^3$ . <sup>b</sup>Transferred from Ref. 8.

<sup>\*</sup>This research has been supported by a grant from the Robert A. Welch Foundation.

<sup>\*\*</sup>Permanent address: Department of Chemistry, University of Oslo, Oslo 3, Norway.

Results and discussion. The force field defined above is very convenient from a practical point of view since force constants can be transferred between similar molecules. For simplicity the force field of methylthiocyanate was calculated first. The results are in good agreement with those previously published by Crowder 2 and Lett and Flygare. The force field of ethylthiocyanate is listed in Table 1. The force constants are quite similar to those of methylthiocyanate and of thioalkanes.8 The calculated and observed frequencies are listed in Table 2 together with an approximate description of the vibrations. With a few exceptions, the assignments are the same as reported in Ref. (4). The bands

Table 2. Observed and calculated frequencies and approximate description of vibratons in CH<sub>3</sub>CH<sub>2</sub>SCN.

${ m Observed}^a \ ({ m cm}^{-1})$	$_{(\mathrm{cm}^{-1})}^{\mathrm{Calculated}}$	Approximate description
2999	2999.1	CH <sub>3</sub> asym. str.
2984	2988.7	$CH_3$ asym. str.
2948	2961.2	CH <sub>2</sub> asym. str.
$2946^{b}$	2949.5	CH <sub>2</sub> sym. str.
2888	2880.3	CH <sub>3</sub> sym. str.
2170	2170.1	CN str.
1456	1463.1	CH <sub>3</sub> scissor
1456	1451.4	CH <sub>3</sub> scissor
1434	1436.9	CH <sub>2</sub> scissor
1389	1386.2	CH <sub>3</sub> sym. def.
1283	1272.6	CH, wag
$1244^b$	1248.1	CH, twist
1066	1058.2	CH <sub>3</sub> rock
$1049^{b}$	1046.2	CH, rock
971	971.8	CC str.
774	793.6	CH, rock
675	685.4	CS str.
640	640.8	$SC \equiv str.$
453 <sup>c</sup>	453.0	SCN bend
394	393.8	SCN bend
326	323.6	CCS bend
$249^b$	249.8	$CH_s$ torsion
149	149.9	CSČ bend
_	$72.2^d$	C-S torsion

<sup>a</sup>Observed frequencies taken from Ref. 4. Unless marked, IR vapour phase values. bLiquid IR value. Liquid Raman value. dThis calculated frequency depends mainly on the transferred force constant value for  $Y_{SC}$ .

at 453 and 394 cm<sup>-1</sup> were assigned to C-C-S bend and the bands at 326 and 308 cm<sup>-1</sup> as S -C≡N bend. The calculations clearly indicate that the S-C≡N bending vibrations are the ones previously assigned as C-C-S bend (453 and 394 cm<sup>-1</sup>) and that the C-C-S bend is at

326 cm<sup>-1</sup>. This is more in accordance with the assignment of the corresponding vibrations in methylthiocyanate, where the two  $S-C\!\equiv\!N$ bending vibrations are 460 and 389 cm<sup>-1.9,10</sup> In CH<sub>3</sub>CH<sub>2</sub>SCH<sub>3</sub> 11 the CCS vibration is observed at 335 cm<sup>-1</sup>, which also is in agreement with the reassignment of these vibrations. Using the force constants found for the C-S torsion reported for several thioalkanes,6 this vibration was calculated to be 72 cm<sup>-1</sup>, in accordance with the value calculated by Ohsaku et al.11 for the C-S torsion in CH<sub>3</sub>ČH<sub>2</sub>-SCH<sub>3</sub>. Only one vibration was observed below 200 cm<sup>-1</sup> the CSC bend at 149 cm<sup>-1</sup>. Although it is not very likely, it is possible that there is an overlap of the C-S torsion and the CSC bending at this frequency. The possibility also exists that, despite the thorough studies both in infrafred and Raman, the lowest fundamental vibration is still not observed.

Acknowledgement. Professor J. E. Boggs and his coworkers at The Center for Structural Studies are thanked for several helpful discussions during the course of this work. Stan McNeme is thanked for doing some of the calculations in this work.

- 1. Hirschmann, R. P., Knisley, R. N. and Fassel, V. A. Spectrochim. Acta. 20 (1964)
- 2. Crowder, G. A. J. Mol. Struct. 7 (1971) 147. 3. Bjørseth, A. and Marstokk, K.-M. J. Mol. Struct. 11 (1972) 15.
- 4. Ellestad, O. H. and Torgrimsen, T. J. Mol. Struct. 12 (1972) 79.
- Gwinn, W. D. J. Chem. Phys. 55 (1971) 477.
- Shimanouchi, T. In Henderson, D., Ed., *Physical Chemistry, An Advanced Treatise*, Academic, New York 1970, Vol. IV,
   Chapter 6.
- 7. Lett, R. G. and Flygare, W. H. J. Chem.
- Phys. 47 (1967) 4730.S. Scott, D. W. and El-Sabban, M. Z. J. Mol. Spectrosc. 30 (1969) 317.
- 9. Moritz, A. G. Spectrochim. Acta. 22 (1966)
- 10. Crowder, G. A. J. Mol. Spectrosc. 23 (1967) 108.
- 11. Ohsaku, M., Shiro, Y. and Murata, H. Bull. Chem. Soc. Jap. 46 (1973) 1399.

Received November 15, 1973.