The 125Te Mössbauer Effect in FeTe2, CoTe2, and RuTe2

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 $^{125}{\rm Te}$ Mössbauer spectra of FeTe2, CoTe2, and RuTe2 have been obtained at 4.2 K. These results, together with the literature data for MnTe2, are discussed in terms of the structural properties of the compounds.

Mössbauer spectroscopy, using the 14.4 $(3/2 \rightarrow 1/2)$ and 37.2 keV $(7/2 \rightarrow 5/2)$ transitions in ⁵⁷Fe and ¹²¹Sb, respectively, has previously been employed by the present authors 1,2 in studying the chemical bonding in (TX_2) compounds with the pyrite (FeS₂-p), marcasite (FeS₂-m), and arsenopyrite (FeAsS; binary prototype CoSb₂) type crystal structures. We here report on the ¹²⁵Te Mössbauer effect (viz. the 35.6 keV $(3/2 \rightarrow 1/2)$ transition) in the compounds FeTe₂, CoTe₃, and RuTe₂ which take the FeS₂-m or FeS₂-p type structure.

EXPERIMENTAL

The stoichiometric FeTe₂ and CoTe₂ samples were prepared and characterized by Guinier photographic X-ray data as described previously. A RuTe₂ sample was prepared by heating a stoichiometric mixture of the elements in an evacuated, sealed quartz tube a 550°C for 14 days. The Guinier data of the product gave a = 6.3903(3) Å, in excellent agreement with the recent literature value of a = 6.3906(4) Å.⁴

The ¹²⁵Te Mössbauer spectra were obtained at (P.C.M.U.,) Harwell, United Kingdom. The source (Cu^{125} Te, viz. ¹²⁵I annealed into Cu) and absorber were maintained at 4.2 K.

The preparation of absorbers and other experimental details, including the computational work, were performed as described previously.2

RESULTS AND DISCUSSION

Fig. 1 shows the ¹²⁵Te Mössbauer spectra of FeTe₂, CoTe₂, and RuTe₂ at 4.2 K. It can be seen that the chemical shifts (δ) and total line-widths of these spectra are equal within experimental error, although their maximum absorption effects differ. There are no resolved quadrupole or magnetic hyperfine interactions observed for any of the compounds. However, the peaks appear to be too broad to be accounted for by single lines and consequently they were computer fitted assuming quadrupole split lines of equal width and intensity.

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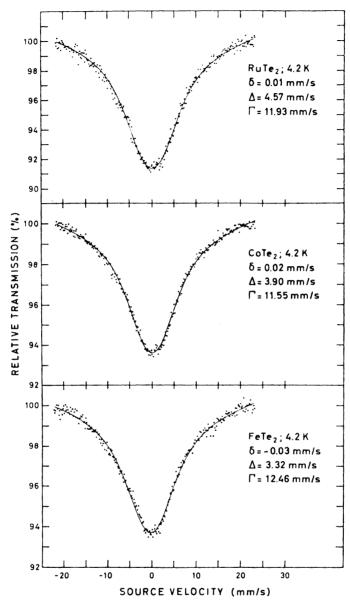


Fig. 1. 125 Te Mössbauer spectra for FeTe₂, CoTe₂, and RuTe₂ at 4.2 K. The curves show the computer fitting to the experimental points.

The resulting values for the parameters δ , $\Delta(=\frac{1}{2}|eQV_{zz}|)$, and Γ are indicated on the diagram. RuTe₂ is diamagnetic,⁵ which immediately excludes the possibility of a magnetic exchange interaction in this compound. The almost

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identical appearance of the three spectra (Fig. 1) consequently suggests that magnetic exchange interactions are also absent in FeTe₂ and CoTe₂ at 4.2 K. Certain indications from magnetic susceptibility data ^{6–10} appear to suggest cooperative magnetic phenomena in FeTe₂ and CoTe₂ and this has been positively claimed in the case of FeTe₂ which is said to be antiferromagnetic below a Néel temperature of 83–85 K.^{7,8} On the other hand, the present finding concerning FeTe₂ is consistent with its ⁵⁷Fe Mössbauer spectrum at 81 K in which a magnetic hyperfine interaction is not apparent.¹¹

It is interesting to compare these observations with those reported by Pasternak and Spijkervet ¹² for MnTe₂ (FeS₂-p type structure). Their ¹²⁵Te Mössbauer spectra above and below the Néel temperature of 83.8 K exhibit resolved quadrupole splittings giving a temperature independent value $eQV_{zz} = -15.5$ mm/s for the quadrupole coupling constant. In addition, the spectra recorded for the antiferromagnetic state of MnTe₂ show significant changes in spectral shape and line width from those recorded for the paramagnetic state. The internal magnetic field induced at the Te nuclei takes a constant value of 114 ± 7 kOe between 4.2 and 20.3 K, decreasing to 55 ± 3 kOe at 77.3 K. None of these features are detected in the spectra of FeTe₂, CoTe₂, and RuTe₂.

The quadrupole coupling which is present in FeTe₂, CoTe₂, and RuTe₂, as a consequence of the distortion from T_d symmetry of the coordination polyhedra around Te, is only resolved by resorting to computer fitting. A measure of the angular part of the distortion from cubic symmetry is given by the average deviation (ξ) of the bond angles from the tetrahedral value of 109.5°, as described in a preceding paper.2 The values of ξ (ranging from 8.3° for AuSb₂ to 13.6° for FeSb₂) have been found to correlate well with the quadrupole coupling constants obtained from the 121 Sb Mössbauer spectra of TSb, compounds with the FeS₂-p, FeS₂-m, and CoSb₂ type structures.² (The ξ values for MnTe₂, FeTe₂, CoTe₂, and RuTe₂ are 6.3, 10.1, 9.9, and 8.9°, respectively.) A corresponding relationship between ξ and eQV_{zz} should also be valid among the ditellurides provided the bonding characteristics of the T atoms are similar within the series under consideration. In MnTe₂ there is a high spin d^5 $(t_{2g}^{3}e_g^{*2})$ configuration associated with the Mn atoms.¹³ This gives rise to longer and weaker Te-Mn bonds and shorter, stronger Te-Te bonds. Hence, the values for ξ and eQV_{zz} for this compound are likely to deviate from the relationship which may be devised for the corresponding parameters of the (low spin) compounds FeTe₂, CoTe₂, and RuTe₂. The fact that, for the TTe₂ compounds, there does appear to be some correlation between ξ and $|\Delta|$ may well be fortuitous in view of the undetermined sign of Δ for FeTe₂, CoTe₂, and RuTe₂.

The positive sign for V_{zz} in MnTe₂ distinguishes this compound from the ten $T\mathrm{Sb}_2$ compounds discussed previously,² which consistently have a negative sign for V_{zz} . In the latter series of compounds the negative sign has been interpreted on the basis of an anisotropic expansion of the Sb 5p orbitals into the bonding orbitals. This implies that there is a net excess of electron density localized in the Sb $5p_z$ orbitals, the z direction defined as being along the X-X pair, cf. Ref. 2. For MnTe₂ where the Te-Mn bonds are weakened the major residual 5p population is localized in the $5p_x$ and $5p_y$ Te atomic orbitals. The smaller quadrupole interactions in FeTe₂, CoTe₂, and RuTe₂ are consistent with this interpretation since an increase in the delocalization of the $5p_x$ and

 $5p_v$ electron density into the (stronger) bonding Te – T orbitals would reduce the quadrupole coupling constant compared with the value observed for MnTe₂.

An analysis of the structural data for compounds with the FeS₂-p, FeS₂-m, and CoSb₂ type structures suggests that the bonding in these compounds is of a predominantly covalent nature. The T-X bond lengths in the high spin manganese dichalcogenides are considerably longer than those for the other 3d transition metal dichalcogenides. As mentioned above the elongation of the Mn-X bond length is attributed to the influence of the half filled e_e * orbitals. Since this is a crucial point of the model it is gratifying to note that this view is supported by the recent ⁵⁷Fe Mössbauer data of Bargeron et al. ¹⁴ These authors substituted 2 % 57Fe into samples of MnS₂, MnSe₂, and MnTe₂ and found that at atmospheric pressure the quadrupole interaction is characteristic of high spin Fe. This should be compared with the situation which prevails in the low spin compounds FeS₂, FeSe₂, and FeTe₂. There is accordingly a definite correspondence between the T-X bond length and e_r * population. In fact, the high spin situation which prevails in MnX_2 can be converted to low spin by squeezing the atoms together under the application of high pressure.

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