singlets (Fig. 1). In cases where this signal could not be resolved at 60 Mc/s (hemiacetals of less sterically hindered alcohols such as 2-butanol) it was possible to resolve this signal into two components by use of 100 Mc/s spectra. The occurrence of epimerization was demonstrated by recording the NMR-spectrum of the previously mentioned hemiacetal of 3,3-dimethyl-2butanol as a function of time (Fig. 2). As shown in the figure, the mole fraction  $X_A$ of the most shielded H<sub>x</sub>-proton changed from 0.72 to 0.50 in 1440 min. No further change in the mole fraction was observed after 1440 min, in contrast to the reaction followed by UV which was complete after 10 min. The mole fraction of the two diastereomers was determined by measuring (by use of a planimeter) the areas of the two signals. To make this measurement possible the signals were separated from each other by drawing a vertical line from the crossing point between the two signals to the base line. It is interesting to observe (Fig. 1) that the mole fraction at equilibrium for three different hemiacetals is approximately 0.5. This means that the free energy difference between the two

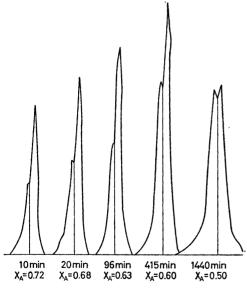


Fig. 2. Time dependence of the NMR-signal at 4.0  $\tau$  of the chloral hemiacetal of 3,3-dimethyl-2-butanol. These spectra were recorded with a xweep-width of 50 cps. 1 mm= $8.35\times10^{-3}$   $\tau$ .  $X_{\rm A}$  is the mole fraction of the most shielded diastereomer.

diastereomeric hemiacetals is approximately zero, contrary to the energy difference between their corresponding transition states. In order to obtain a better understanding of the epimerization process, a study of the half lives of various hemiacetals as a function of the steric bulk of the groups surrounding the asymmetric centre is planned.

Experimental. NMR Spectra were recorded on a Varian A-60 A spectrophotometer in  $CCl_4$  solution with TMS as internal standard at 25° The concentration of chloral and the alcohols were both  $5\times 10^{-1}$  M and the concentration of acetic acid was  $10^{-2}$  M in the experiments shown in Fig. 1. For the experiments shown in Fig. 2, no acetic acid was added and the concentrations of chloral and 3,3-dimethyl-2-butanol were both  $5\times 10^{-1}$  M.

Acknowledgement. The author is much indebted to Professor K. A. Jensen for his kind interest and helpful criticism.

- Jensen, R. B. Acta Chem. Scand. 17 (1963) 1463.
- Jensen, R. B. and Munksgaard, E. C. Acta Chem. Scand. 18 (1964) 1896.
- Jensen, R. B. and Butty, M. Acta Chem. Scand. 22 (1968) 477.
- Jensen, R. B. Acta Chem. Scand. 22 (1968) 1271.
- Hudson, C. S. Z. physik. Chem. (Leipzig) 44 (1903) 487.

Received July 11, 1968.

## An SCF MO Study of Two Isomeric Thiadiazoles

PETER MARKOV\* and P. N. SKANCKE

Department of Chemistry, University of Oslo, Blindern, Oslo 3, Norway

From a theoretical point a view the electronic structure of the monocyclic thiadiazoles is of considerable interest.<sup>1-3</sup> Previous calculations, performed within the Hückel-approximation, have made possible a discussion of some of the molec-

<sup>\*</sup> Permanent address: Department of Organic Chemistry, University of Sofia, Bulgaria.

ular properties related to the  $\pi$ -electron distributions.<sup>2,3</sup>

The purpose of the present note is to present a comparative study of the π-electron structures of the ground states of the symmetrical isomers 1,3,4- and 1,2,5-thiadiazole, called (I) and (II), respectively, by using the well-known Pariser-Parr-Pople approximation 4-6 in the SCF theory.

Our calculations are based on the assumption that the sulphur atom contributes a  $3p\pi$ -orbital to the conjugate system.

The semi-empirical parameters have been evaluated in the following way:

The core resonance integrals,  $\beta_{\mu\nu}$ , were determined by use of the simple relation

$$\beta \mu \nu = \beta_0 S \mu \nu / S \tag{1}$$

where  $S_0$  and  $\beta_0$  are the overlap and resonance integrals for benzene. The latter was assumed to be  $-2.39 \, \mathrm{eV}^{.5,7}$  The evaluation of overlap integrals was based on Slater type orbitals, and experimental bond distances. 14,16

The diagonal elements of the core operator were decomposed according to the method of Goeppert-Mayer and Sklar:<sup>8</sup>

$$\alpha \mu = W \mu - \sum_{\nu \neq \mu} n^{\nu} \gamma_{\mu\nu} \tag{2}$$

In (2)  $n\nu$  is the number of  $\pi$ -electrons contributed to the system by atom  $\nu$ .  $\gamma\mu\nu$  is the Coulomb repulsion integral between electrons centered at atoms  $\mu$  and  $\nu$ , and  $W_{\mu}$  may be interpreted as an approximation to the valence state ionization potential for atom  $\mu$ . In our calculations we have used the values  $^{0}$   $W_{\rm C} = -11.54$  eV,  $W_{\rm N} = -14.32$  eV, and  $W_{\rm S} = -12.50$  eV.

These values are not quite appropriate for the valence states considered relevant here, but we found it to be an advantage to have W-values referring to corresponding states for all the atoms. Furthermore, the changes in the values due to rearrangements of the underlying  $\sigma$ -system are very small.<sup>10</sup>

The one-center two-electron integrals were evaluated by use of the relation <sup>5</sup>

$$\gamma_{\mathbf{X}}^{\mathbf{0}} = (\mathbf{IP})_{\mathbf{X}} - (\mathbf{EA})_{\mathbf{X}} \tag{3}$$

where the atomic IP-values are those referred to above, and EA represents the electron affinity for the corresponding states. The values obtained are  $\gamma_{\rm S}^0=9.80$  eV,  $\gamma_{\rm C}^0=10.53$  eV, and  $\gamma_{\rm N}^0=12.73$  eV. The two-center two-electron integrals were estimated by means of the uniformly charged sphere approximation, with orbital exponents based on Slaters rules.

The bond orders and  $\pi$ -electron charges obtained by the calculations, which were performed on a UNIVAC 3300 computer, are collected in Table 1 and Table 2, respectively. For numbering of atoms see

Fig. 1.

The calculated bond orders were in the usual way employed for a prediction of bond distances. For the C—C bond in (II) we used the empirical linear relation

$$R_{\rm CC} = 1.517 - 0.18 p_{\rm CC} \tag{4}$$

discussed previously by one of us.  $^{11}$  A similar relation for C-N bonds has recently been suggested  $^{12}$ 

$$R_{\rm CN} = 1.458 - 0.18 p_{\rm CN} \tag{5}$$

Table 1. Calculated bond orders for 1,3,4-thiadiazole (I) and 1,2,5-thiadiazole (II). Values obtained by Hückel calculations (HMO) are included for comparison.

|      | Molecule<br>I | Molecule<br>II       |  |
|------|---------------|----------------------|--|
| Bond | SCF HMOª      | SCF HMO              |  |
| 1-2  | 0.3392 0.464  | 0.3109 0.468a 0.536b |  |
| 2-3  | 0.8902 0.787  | 0.7452 0.780 0.676   |  |
| 3-4  | 0.3099 0.526  | 0.3830 0.541 0.680   |  |

<sup>&</sup>lt;sup>a</sup> Ref. 2; <sup>b</sup> Ref. 3.

Table 2. Calculated π-electron charges on the different atoms. In electron units.

| Atom | Molecule<br>I | Molecule<br>II |  |
|------|---------------|----------------|--|
| 1    | +0.1647 (S)   | +0.2656 (S)    |  |
| 2    | +0.1702 (C)   | -0.5227 (N)    |  |
| 3    | -0.2525 (N)   | +0.3899 (C)    |  |

Fig. 1. Numbering of atoms in 1,3,4-thiadiazole (I) and 1,2,5-thiadiazole (II).

Acta Chem. Scand. 22 (1968) No. 6

By using the calculated bond order and the measured N-N distance in pyridazine, we obtain the following relation 13

$$R_{\rm NN} = 1.424 - 0.18 p_{\rm NN} \tag{6}$$

Application of these formulae gives the bond distances presented in Table 3, where also experimental values obtained by microwave spectroscopy 14,15 are included for comparison.

Table 3. Calculated and observed bond distances. In Å units.

| Bond | Molecule Molec  |                   | ıle               |
|------|---|-------------------|-------------------|
|      | $egin{array}{c} \mathbf{I} \ R_{ m calc.} & R_{ m obs.} ^{a} \end{array}$ | $R_{ m calc.}$ II | $R_{ m obs.}{}^a$ |
| 1-2  | - 1.7205  | _                 | 1.631             |
| 2-3  | 1.298 1.3024  | 1.324             | 1.328             |
| 3-4  | 1.368 1.3710  | 1.448             | 1.420             |

Observed R-values: a Ref. 14; b Ref. 15.

The values in Table 1 clearly show that the low bond orders obtained by the Hückel approximation are decreased by the SCF calculation, and the high bond orders are correspondingly increased. This accentuation of alternating high and low  $\pi$ electron densities is rather satisfactory due to the presence of hetero-atoms with different electronegativities. The agreement between calculated and observed distances shown in Table 3 is very satisfactory for all the bonds considered except for the C-C bond in (II). Also here the agreement has to be considered as acceptable. Particularly the observed difference between the N-C bond distances in the two isomers is exactly reproduced by our calculations.

The predicted vertical ionization potentials are 12.26 eV and 10.00 eV for (I) and (II), respectively. These values are most certainly too high, due to the fact that the applied  $W_{\mu}$  values are based on atomic valence state energies. However, the calculated difference should be rather realistic. To the authors' knowledge no experimental values are available.

The calculated atomic charges presented in Table 2 indicate a high degree of polarity in both isomers. The predicted π-electron dipole moments are 3.86 D and 2.92 D for (I) and (II), respectively. The only observed value, <sup>16</sup> 3.28 D for (I), is qualitatively in agreement with the calculated value. The values obtained for the π-electron charge distributions permit us to conclude that (II) should be more sensitive to both nucleophilic and electrophilic attacks than (I). This conclusion is in accordance with available experimental information.<sup>1,2</sup>

- Elderfield, R. C. Heterocyclic Compounds, Wiley, New York 1951, Vol. 7, p. 585.
- Zahradnik, R. and Koutecky, J. Collection Czech. Chem. Commun. 26 (1961) 156.
- Ray, N. K. and Narasimhan, P. T. Theoret. Chim. Acta 5 (1966) 401.
- Pariser, R. and Parr, R. G. J. Chem. Phys. 21 (1953) 466.
- Pariser, R. and Parr, R. G. J. Chem. Phys. 21 (1953) 767.
- Pople, J. A. Trans. Faraday Soc. 49 (1953) 1375.
- Fischer-Hjalmars, I. In Löwdin, P. O. and Pullman, B. Molecular Orbitals in Chemistry, Physics and Biology, Academic, N.Y. 1964.
- Goeppert-Mayer, M. and Sklar, A. L. J. Chem. Phys. 6 (1938) 645.
- Skinner, H. A. and Pritchard, H. O. Trans. Faraday Soc. 49 (1953) 1254.
- Pilcher, G. and Skinner, H. A. J. Inorg. Nucl. Chem. 24 (1962) 937.
- Skancke, P. N. Acta Chem. Scand. 18 (1964) 1671.
- Fischer-Hjalmars, I. and Sundbom, M. Acta Chem. Scand. 22 (1968) 607.
- Gropen, O. and Skancke, P. N. To be published.
- Bak, B., Nygaard, L., Pedersen, E. J. and Rastrup-Andersen, J. J. Mol. Spectry. 19 (1966) 283.
- Dobyns, S. V. and Pierce, L. J. Am. Chem. Soc. 85 (1963) 3553.
- Bak, B., Christensen, D., Nygaard, L., Lipschitz, L. and Rastrup-Andersen, J. J. Mol. Spectry. 9 (1962) 225.

Received June 27, 1968.