The Crystal Structure of Strontium Diborate, SrO · 2B₂O₃

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Strontium diborate, ${\rm Sr0\cdot 2B_2O_3}$, crystallizes in the space group $Pmn2_1$ with unit cell dimensions: $a=10.711\pm0.005$ Å, $b=4.427\pm0.002$ Å, $c=4.235\pm0.002$ Å and two formula units in the cell, corresponding to a calculated density of 4.02 g/cm³. The crystal structure has been solved with two-dimensional Patterson and three-dimensional Fourier syntheses and refined with the method of least squares. A reliability index R=0.126 was obtained for data from 189 three-dimensional reciprocal lattice points. The structure consists of a three-dimensional network of BO₄-tetrahedra. Some of the oxygens in the structure are bonded to three borons.

As part of a general survey of anhydrous polyborate structures, the crystal structure of strontium diborate has been determined. The preceding paper in this series is concerned with the structure of silver tetraborate. This compound obeyed the rule that the number of borons in four-fold coordination is equal to the number of foreign cation charges. The same coordination behaviour of boron was previously observed in a series of alkali borates with a high boron oxide content. Considering the small energy requirement for a coordination change of boron, it cannot be assumed that the rule is of general validity, especially if other cations than alkali atoms are introduced. The present investigation has revealed a new case where the rule breaks down. Here we have a type of anomalous coordination behaviour related to that discovered by Clark in the hydrated borate mineral, tunnellite, $SrO\cdot3B_2O_3\cdot4H_2O$.

Recently X-ray data for strontium diborate have been published by Stewart and Rindone⁴, and by Block, Perloff and Weir.⁵ The unit cell dimensions given by the latter authors agree reasonably well with those reported here.

EXPERIMENTAL

Strontium diborate was crystallized around 900°C from a melt with a slight excess of boron oxide.

The following unit cell dimensions were determined with a least squares fit of powder diffraction data from a Hägg-Guinier camera, using lead nitrate as internal standard:

$$a = 10.711 \pm 0.005 \text{ Å}$$
 $b = 4.427 \pm 0.002 \text{ Å}$
 $c = 4.235 \pm 0.002 \text{ Å}$

With 2 formula units of SrO:2B₂O₃ in the cell, the calculated density becomes 4.02 g/cm³. This is a relatively large value. However, by floating the sample in Clericis solution it

was verified that the density was above 3.9 g/cm³ and close to 4.0 g/cm.³

Single crystal integrated Weissenberg exposures were obtained with $CuK\alpha$ -radiation around the two shortest axes for the zeroth to second layer. Systematic extinction were observed for the h0l reflections when h+l=2n+1, indicating one of the space groups $Pmn2_1$ and Pmmn.

The crystal, which was rotated around the b-axis, had a quadratic cross-section 0.005×0.005 cm perpendicular to the axis of rotation and the crystal rotated around the c-axis had a corresponding cross-section 0.007×0.0085 cm. The linear absorption coefficient is 200 cm^{-1} for $\text{Cu}K\alpha$ -radiation, so that the absorption effects are not entirely negligible. No corrections for absorption were made, but a rough estimate of the effect

will be given in the next paragraph.

Intensities were estimated visually from triple film packs by comparing the Weissenberg spots with a calibrated strip. After the Lorentz-polarization corrections, the intensities from the various layers around the two axes were adjusted so as to bring all the values on a common scale. For reciprocal lattice points, where two measurements were available, those from the smaller crystal were given a higher weight.

THE STRUCTURE DETERMINATION

The Patterson projections revealed a strontium—strontium vector incompatible with the special two-fold positions in the space group Pmmn. It was therefore concluded that the correct space group is $Pmn2_1$. Based on information from the Patterson functions about the strontium position, a three-dimensional electron density distribution was calculated. The summation contained a false symmetry (since phases were obtained only from strontium positions) but peaks due to the false symmetry could be eliminated by considering interatomic distances. The structure thus obtained was refined to a reliability index of 12.6 % with the method of least squares, using the program written by Åsbrink and Brändén for the Facit computer. Absolute values of the calculated structure factors are given in Table 1 and the final atomic parameters are reported in Table 2. The atoms Sr and O_1 occupy special two-fold positions in the space group. The origin along the z-axis is arbitrarily fixed on the strontium atom.

From the two last columns given in Table 2, the strontium atom appears to have a significant negative temperature factor parameter B=-0.67. This is largely accounted for by absorption errors, however. If the crystal with cross-section 0.005×0.005 cm (respectively 0.007×0.0085 cm) had been a cylinder with the same cross-sectional area, the absorption correction for the equatorial layer would, well within the accuracy of the intensity reading, correspond to multiplying the intensities with an artificial temperature factor with parameter value 2B=0.4 (respectively 2B=0.6)... Somewhat larger corrections are valid for the upper layers, and larger corrections also apply for certain planes in the equatorial layer of the quadratic or rectangular crystals. It is therefore reasonable to blame the absorption errors for the fact that the temperature factor parameter for strontium becomes negative in the least squares refinement.

Table 1. Observed and calculated structure factors, multiplied by 10.

h k l	$F_{\mathbf{o}}$	$ m{F_c} $	h k l	$F_{\mathbf{o}}$	$ m{F_c} $	h k l	$oldsymbol{F_{\mathbf{o}}}$	$ F_{\mathbf{c}} $
2 0 0	263	328	3 1 2	469	414	8 2 2	566	483
400	422	469	4 1 2	384	226	9 2 2	302	296
600	601	625	5 1 2	530	482	10 2 2	363	395
$\begin{smallmatrix}8&0&0\\10&0&0\end{smallmatrix}$	$\frac{910}{352}$	$\begin{array}{c} 914 \\ 333 \end{array}$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 106 \\ 501 \end{array}$	$\begin{array}{c} 92 \\ 492 \end{array}$	$\begin{smallmatrix}0&2&3\\1&2&3\end{smallmatrix}$	$255 \\ 331$	165
$10 \ 0 \ 0$	310	325	$\begin{smallmatrix}7&1&2\\8&1&2\end{smallmatrix}$	79	90	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	366	301 375
101	404	323 377	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	390	396	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	575	494
3 0 1	645	765	10 1 2	129	112	$\begin{smallmatrix}3&2&3\\4&2&3\end{smallmatrix}$	134	182
5 0 1	475	548	11 1 2	351	368	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	584	459
7 0 1	504	467	0 1 3	554	538	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	343	319
$9 \tilde{0} \tilde{1}$	366	370	1 1 3	255	192	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	231	267
11 0 1	525	527	2 1 3	736	646	8 2 3	137	162
0 0 2	889	858	4 1 3	457	287	$0\ 2\ 4$	334	416
202	563	653	5 1 3	128	134	124	196	255
402	489	517	6 1 3	513	568	$2\ 2\ 4$	275	396
$6 \ 0 \ 2$	519	549	7 1 3	116	112	3 2 4	160	208
802	621	650	8 1 3	443	510	424	434	428
10 0 2	343	392	9 1 3	181	189	0 3 0	316	283
103	396	332	114	557	427	1 3 0	246	214
303	581	593	214	149	125	230	272	225
503	581	655	3 1 4	351	409	3 3 0	472	434
7 0 3	119	187	414	114	135	4 3 0	692	728
9 0 3	419	455	5 1 4	466	465	5 3 0	425	437
0 0 4	401	488	6 1 4	81	123	6 3 0	152	165
2 0 4	478	506	0 2 0	778	985	7 3 0	90	120
4 0 4	390	473	1 2 0	178	180	8 3 0	248	219
604	325	419	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	434	505	930	325	337
$\begin{smallmatrix}0&1&0\\1&1&0\end{smallmatrix}$	$\begin{array}{c} 190 \\ 272 \end{array}$	161	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 434 \\ 422 \end{array}$	414	$\begin{smallmatrix}0&3&1\\1&3&1\end{smallmatrix}$	460	431
$\begin{array}{cccc} 1 & 1 & 0 \\ 2 & 1 & 0 \end{array}$	44	351	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	422 76	377 75	$\begin{smallmatrix}1&3&1\\2&3&1\end{smallmatrix}$	$\begin{array}{c} 504 \\ 607 \end{array}$	$\begin{array}{c} 428 \\ 472 \end{array}$
3 1 0	50 4	17 689	$\begin{smallmatrix} 5 & 2 & 0 \\ 6 & 2 & 0 \end{smallmatrix}$	583	518	$\begin{smallmatrix}2&3&1\\3&3&1\end{smallmatrix}$	284	232
4 1 0	487	518	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	334	277	431	451	367
5 1 0	451	450	8 2 0	675	636	5 3 1	213	206
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	49	12	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	166	165	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	484	413
7 1 0	510	506	$10 \begin{array}{cccccccccccccccccccccccccccccccccccc$	369	344	7 3 1	484	431
8 1 0	59	62	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	316	330	8 3 1	419	367
9 1 0	255	238	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	580	459	$\stackrel{\circ}{9}\stackrel{\circ}{3}\stackrel{\circ}{1}$	355	309
10 1 0	167	163	$\tilde{1}$ $\bar{2}$ $\bar{1}$	475	423	10 3 1	375	330
11 1 0	525	501	$\overline{2}$ $\overline{2}$ $\overline{1}$	313	254	0 3 2	331	269
12 1 0	161	162	3 2 1	486	433	132	384	317
0 1 1	513	536	4 2 1	211	212	2 3 2	360	344
111	316	259	5 2 1	443	419	3 3 2	451	341
$2\ 1\ 1$	700	784	6 2 1	434	365	$4 \;\; 3 \;\; 2$	598	477
3 1 1	108	131	7 2 1	463	388	5 3 2	407	293
4 1 1	680	783	8 2 1	243	217	6 3 2	266	257
5 1 1	96	69	9 2 1	346	349	7 3 2	372	337
6 1 1	504	508	10 2 1	164	181	8 3 2	296	258
7 1 1	355	247	11 2 1	357	363	9 3 2	313	277
8 1 1	495	421	$\begin{array}{cccc} 0 & 2 & 2 \\ 1 & 2 & 2 \end{array}$	645	578	0 4 0	166	114
9 1 1	119	111	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	266	220	1 4 0	437	414
$\begin{array}{cccc} 10 & 1 & 1 \\ 11 & 1 & 1 \end{array}$	$\begin{array}{c} 656 \\ 93 \end{array}$	$\begin{array}{c} \textbf{642} \\ \textbf{93} \end{array}$	$\begin{smallmatrix}2&2&2\\3&2&2\end{smallmatrix}$	545	$\begin{array}{c} 484 \\ 272 \end{array}$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	434 363	$\begin{array}{c} 424 \\ 333 \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	93 378	93 371	$\begin{smallmatrix}3&2&2\\4&2&2\end{smallmatrix}$	360 49 6	506	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 303 \\ 278 \end{array}$	333 239
$\begin{smallmatrix}12&1&1\\0&1&2\end{smallmatrix}$	$\frac{378}{125}$	187	$\begin{smallmatrix}4&2&2\\5&2&2\end{smallmatrix}$	$\begin{array}{c} 490 \\ 272 \end{array}$	281	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 278 \\ 246 \end{array}$	$\begin{array}{c} 239 \\ 277 \end{array}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	660	605	$\begin{smallmatrix} 5&2&2\\6&2&2\end{smallmatrix}$	507	416	$\begin{array}{c} 540 \\ 640 \end{array}$	331	$\frac{277}{312}$
$\begin{smallmatrix}1&1&2\\2&1&2\end{smallmatrix}$	219	214	$\begin{smallmatrix} 0 & 2 & 2 \\ 7 & 2 & 2 \end{smallmatrix}$	193	227	740	$\frac{331}{337}$	$\begin{array}{c} 312 \\ 402 \end{array}$
4 1 4	410	417	. 4 4	199	441	, 1 U	001	TU 4

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h k l	F_{o}	$ F_{c} $	h k l	$F_{\rm o}$	$ F_{c} $	h k l	F_{o}	$ \boldsymbol{F}_{\mathbf{c}} $
041	484	362	741	360	312	642	334	248
141	437	337	042	210	183	050	293	345
241	522	433	142	437	373	1 5 0	61	83
341	155	101	242	396	282	250	340	433
441	428	349	3 4 2	407	366	0 5 1	193	229
5 4 1	246	202	442	413	414	151	387	401
641	516	470	5 4 2	360	321	2 5 1	222	151

Table 2. Atomic position parameters (with standard deviations) as fractions, multiplied by 1000, of the unit cell edge. The last two columns give the parameter B (together with its standard deviation) of the temperature factor $\exp[-B(\sin\Theta/\lambda)^2]$.

Atom	x/a	σ _{x/a}	y/b	συ/β	z/c	σz c	B (in Å ²)	$\sigma_{\mathtt{B}}$
Sr	0		289	1	0		-0.67	0.07
O ₁ O ₂ O ₃	0		728	12	454	20	1.1	1.1
O,	359	2	857	7	064	8	-0.6	0.5
O_3	221	3	631	8	335	9	0.0	0.6
O,	365	4	226	9	335	12	0.6	0.8
B,	379	4	174	9	976	20	-0.6	0.7
O ₄ B ₁ B ₂	246	5	671	11	963	19	0.1	0.9

DESCRIPTION AND DISCUSSION OF THE STRUCTURE

A projection of the structure along the a-axis is shown in Fig. 1. The structure consists of a three-dimensional network of BO₄-tetrahedra. The

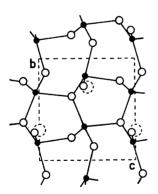


Fig. 1. Projection along the a-axis. Filled circles represent boron, open circles represent oxygen, large dashed circles represent strontium. The unit cell is indicated with dashed lines. Oxygens closest to strontium in projection form bridges with layers below or above, which are mirror images of the layer shown.

layer of BO_4 -tetrahedra, shown in the figure, is connected with its own mirror image exactly above or below through the atom O_1 . The O_1 atom, seen closest to strontium in the projection, will thus appear in this projection as if it were bonded to only one boron. (O_1 atoms with z-parameters 0.454 connect with layers above, those with z-parameters 0.954 connect with layers below).

Table 3. Boron—oxygen bond distances and shortest strontium—oxygen distances, with standard deviations. Each of the tabulated distances from strontium to O₂, O₃, and O₄ occurs twice, so that strontium has 9 oxygen neighbours below 3 Å.

$\begin{array}{c} B_1 - O_1 \\ B_1 - O_2 \\ B_1 - O_3 \\ B_1 - O_4 \end{array}$	$egin{array}{ll} 1.37 & \pm & 0.04 \ \hbox{\AA} \\ 1.47 & \pm & 0.05 \\ 1.50 & \pm & 0.06 \\ 1.55 & \pm & 0.09 \end{array}$	$Sr - O_1$	$\begin{array}{c} \textbf{2.73} \pm 0.07 \text{\AA} \\ \textbf{3.02} \pm 0.07 \\ \textbf{3.14} \pm 0.07 \\ \textbf{3.39} \pm 0.07 \end{array}$
$ B_2 - O_2 B_2 - O_3 B_3 - O_3 $	$egin{array}{l} 1.53\ \pm\ 0.06 \ 1.48\ \pm\ 0.06 \ 1.61\ \pm\ 0.09 \ \end{array}$	$Sr - O_2$	$\begin{array}{c} \textbf{2.47} \pm \textbf{0.03} \\ \textbf{2.90} \pm \textbf{0.03} \end{array}$
$B_2 - O_4$	1.38 ± 0.06	$\mathrm{Sr}-\mathrm{O_3}$	$\begin{array}{c} {f 3.09} \pm 0.03 \\ {f 3.15} \pm 0.03 \end{array}$
		$Sr - O_4$	$\begin{array}{c} \textbf{2.68} \pm 0.04 \\ \textbf{2.79} \pm 0.04 \end{array}$

It is also seen from Fig. 1 that the structure contains six-membered boron-oxygen rings. Such rings are a very common feature in a wide variety of borate structures. The present structure may even be described as chains along the c-axis formed by these six-membered rings sharing edges. The chains are crosslinked with neighbouring chains at regular intervals.

The boron-oxygen bond lengths are given in Table 3. The average value of these bond lengths, 1.49 Å, is slightly larger than the average value of 1.47 Å found for BO_4 -tetrahedra in several structures. The spread of the individual bond lengths about the average value cannot be considered significant (perhaps with the exception of the B_1-O_1 distance). It is not unlikely, however, that the BO_4 -tetrahedra could be somewhat distorted.

The atom O₃ is remarkable inasmuch as it is bonded to three boron atoms. This situation is not quite unique, however. Structures with oxygens three-fold coordinated by boron have been reported for hexagonal boron oxide by Berger ⁶ and recently for the mineral tunellite, SrO·3B₂O₃·4H₂O, by Clark.³ Three-bonding oxygens with two bonds to boron and a third to hydrogen are known from some metaboric acids.⁷ In the present structure the average of the three bond distances to O₃ (1.53 Å) seems to be somewhat larger than the average of the remaining boron-oxygen distances (1.46 Å). The available data are not sufficiently accurate to establish this lengthening with certainty, but a similar lengthening of the boron-oxygen bond for three-fold coordinated oxygen has been indicated by the study of tunellite.³

The coordination of strontium by oxygen is definitely irregular. An irregular cation coordination has previously been observed in other anhydrous borates. It is therefore difficult to establish a meaningful value for the coordination number of the cation in these phases. Table 3 lists values for the 16 shortest strontium-oxygen distances ranging upwards from 2.74 Å to 3.39 Å. The sum of the ionic radii of strontium and oxygen is 2.35 Å. Setting an arbitrary limit at 2.90 Å for the first coordination sphere of strontium, a coordination number of 9 follows. It is interesting to note that the three-fold coordinated oxygen is comparatively far removed (3.09 Å or more) from strontium.

The structure reported here demonstrates the precarious balance of boron between the coordination numbers three and four. This circumstance has previously been emphasized by the author, e.g. it has been proposed that the comparatively low viscosity of molten boron oxide is due to the low energy required for activating states with four-fold coordinated boron and threefold coordinated oxygen.8

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