# Vibrational Mean-Square Amplitude Matrices

## VI. Planar Symmetrical XY<sub>3</sub> Molecules

#### S. J. CYVIN

Institutt for teoretisk kjemi, Norges tekniske høgskole, Trondheim, Norway

The mean-square amplitude quantities for the planar symmetrical XY<sub>3</sub> molecular model are treated theoretically. The L matrix elements (S = LQ), and the symmetrized mean-square amplitude matrix elements at the temperatures T = 0 and 298 °K, are given numerically for the four boron tribalides.

The vibrations of planar symmetrical XY<sub>3</sub> molecules have been subjected to several investigations. For publications containing theoretical considerations on this problem, and involving force constant determinations, references are made to 1-7. For the computations of mean amplitudes of vibration from spectroscopic data, it is referred to 7-9. To the knowledge of the author, no determinations of mean amplitudes of vibration have so far been performed from electron-diffraction, for any molecule of the structure here considered. In the present paper, the mean-square amplitude quantities of planar symmetrical XY<sub>3</sub> molecules are treated, and numerical values for boron tribalides are reported.

#### THEORETICAL TREATMENT

For the in-plane vibrations of the considered molecular model there exist four independent harmonic vibrational constants. The following set of constants consists of mean-square amplitude quantities, based on a set of valence force coordinates.

$$\sigma = \langle r_1^2 \rangle = \langle r_2^2 \rangle = \langle r_3^2 \rangle \tag{1}$$

$$\sigma' = \langle r_1 r_2 \rangle = \langle r_2 r_2 \rangle = \langle r_2 r_1 \rangle \tag{2}$$

$$\tau = R^2 \langle \alpha_1^2 \rangle = R^2 \langle \alpha_2^2 \rangle = R^2 \langle \alpha_2^2 \rangle \tag{3}$$

$$\tau' = R^2 \langle \alpha_1 \alpha_2 \rangle = R^2 \langle \alpha_2 \alpha_3 \rangle = R^2 \langle \alpha_3 \alpha_1 \rangle \tag{4}$$

$$\rho = R\langle r_1 \alpha_1 \rangle = R\langle r_2 \alpha_2 \rangle = R\langle r_3 \alpha_3 \rangle \tag{5}$$

$$\sigma = \langle r_1^2 \rangle = \langle r_2^2 \rangle = \langle r_3^2 \rangle 
\sigma' = \langle r_1 r_2 \rangle = \langle r_2 r_3 \rangle = \langle r_3 r_1 \rangle 
\tau = R^2 \langle \alpha_1^2 \rangle = R^2 \langle \alpha_2^2 \rangle = R^2 \langle \alpha_3^2 \rangle 
\tau' = R^2 \langle \alpha_1 \alpha_2 \rangle = R^2 \langle \alpha_2 \alpha_3 \rangle = R^2 \langle \alpha_3 \alpha_1 \rangle 
\varrho = R \langle r_1 \alpha_1 \rangle = R \langle r_2 \alpha_2 \rangle = R \langle r_3 \alpha_3 \rangle 
\varrho' = R \langle r_1 \alpha_2 \rangle = R \langle r_1 \alpha_3 \rangle = R \langle r_2 \alpha_3 \rangle 
= R \langle r_2 \alpha_1 \rangle = R \langle r_3 \alpha_1 \rangle = R \langle r_3 \alpha_2 \rangle$$
(6)

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In these equations, R is the equilibrium X-Y bond length, r representing the deviations of the interatomic X-Y distances, and a the inter-bond angle deviations, from their equilibrium values. The numbering of the coordinates is the same as that used in previous papers 7,9. From the redundant condition  $\alpha_1 + \alpha_2 + \alpha_3 = 0$  it is found

$$\tau + 2\tau' = 0, \qquad \qquad \varrho + 2\varrho' = 0 \tag{7}$$

for the quantities defined by eqns (3)—(6).

The interatomic distance deviations between the non-bonded atom pairs may be taken into account instead of the angle displacements, and designated  $r_1^*$ ,  $r_2^*$  and  $r_3^*$ . Then the following mean-square amplitude quantities are to be introduced.

$$\tau^* = \langle (r_1^*)^2 \rangle = \langle (r_2^*)^2 \rangle = \langle (r_3^*)^2 \rangle 
\rho^* = \langle r, r, * \rangle = \langle r, r, * \rangle = \langle r, r, * \rangle$$
(8)

$$\tau^* = \langle (r_1^*)^2 \rangle = \langle (r_2^*)^2 \rangle = \langle (r_3^*)^2 \rangle 
\varrho^* = \langle r_1 r_1^* \rangle = \langle r_2 r_2^* \rangle = \langle r_3 r_3^* \rangle 
\varrho'^* = \langle r_1 r_2^* \rangle = \langle r_1 r_3^* \rangle = \langle r_2 r_3^* \rangle 
= \langle r_2 r_1^* \rangle = \langle r_3 r_1^* \rangle = \langle r_3 r_2^* \rangle$$
(8)
(9)

These quantities may be expressed in terms of the quantities from eqns. (1)— (6) by the following linear combinations.

$$\tau^* = \frac{3}{2}(\sigma + \sigma') + \frac{1}{6}(\tau - \tau') - 3^{-\frac{1}{2}}(\varrho - \varrho') \tag{11}$$

$$\varrho^* = 3^{\frac{1}{3}} \sigma' + \frac{1}{3} (\varrho - \varrho') \tag{12}$$

$$\varrho'^* = \frac{1}{2} \, 3^{\frac{1}{2}} \, (\sigma + \sigma') - \frac{1}{6} (\varrho - \varrho') \tag{13}$$

The vibrational analysis is facilitated by introducing symmetry coordinates. The elements of the symmetrized mean-square amplitude matrix are given by

$$\Sigma_{1} = \langle S_{1}^{2} \rangle = \sigma + 2\sigma' \tag{14}$$

$$\Sigma_{3} = \langle S_{3a}^{2} \rangle = \langle S_{3b}^{2} \rangle = \sigma - \sigma' \tag{15}$$

$$\Sigma_{4} = \langle S_{4a}^{2} \rangle = \langle S_{4b}^{2} \rangle = \tau - \tau' \tag{16}$$

$$\Sigma_{34} = \langle S_{3a}S_{4a} \rangle = \langle S_{3b}S_{4b} \rangle = \varrho - \varrho' \tag{17}$$

$$\Sigma_3 = \langle S_{3a}^2 \rangle = \langle S_{3b}^2 \rangle = \sigma - \sigma' \tag{15}$$

$$\Sigma_{\mathbf{A}} = \langle S_{\mathbf{4a}^2} \rangle = \langle S_{\mathbf{4b}^2} \rangle = \tau - \tau' \tag{16}$$

$$\Sigma_{34} = \langle S_{3a}S_{4a} \rangle = \langle S_{3b}S_{4b} \rangle = \rho - \rho' \tag{17}$$

Table 1. Symmetry coordinates for the planar symmetrical XY<sub>3</sub> molecular model.

Species	No.	Coordinate
 A <sub>1</sub> ′	1	$3-\frac{1}{5}$ $(r_1+r_2+r_3)$
$\mathbf{E'_a}$	3a 4a	$\begin{array}{cccc} 3^{-\frac{1}{2}} & (r_1 + r_2 + r_3) \\ 6^{-\frac{1}{2}} & (2r_1 - r_2 - r_3) \\ 6^{-\frac{1}{2}} & R & (2a_1 - a_2 - a_3) \\ 2^{-\frac{1}{2}} & (r_2 - r_3) \end{array}$
${f E'_b}$	$ \begin{array}{c} \mathbf{3b} \\ \mathbf{4b} \end{array} $	$\begin{array}{cccc} 0 & 11 & (2a_1 - a_2 & a_3) \\ 2^{-\frac{1}{2}} & (r_2 - r_3) \\ 2^{-\frac{1}{2}} & R(a_2 - a_2) \end{array}$

The symmetry coordinates (S) are formed as normalized linear combinations of the valence force coordinates, and specified in Table 1. For the meansquare amplitude quantities of eqns. (14)—(17) one has the relations

$$\Delta_1 = \Sigma_1 \mu_{\mathbf{Y}}^{-1} \tag{18}$$

$$\Delta_{1} = \Sigma_{1} \mu_{Y}^{-1} \tag{18}$$

$$\Delta_{3} + \Delta_{4} = \left[ \frac{1}{2} (\Sigma_{3} + \frac{1}{3} \Sigma_{4}) (3\mu_{X} + 2\mu_{Y}) - 3^{\frac{1}{2}} \Sigma_{34} \mu_{X} \right] (3\mu_{X} + \mu_{Y})^{-1} \mu_{Y}^{-1} \tag{19}$$

$$\Delta_3 \Delta_4 = \frac{1}{3} (\Sigma_3 \Sigma_4 - \Sigma_{34}^2) (3\mu_X + \mu_Y)^{-1} \mu_Y^{-1}$$
 (20)

Here  $\mu_X$  and  $\mu_Y$  denote the inverse masses of the X and Y atoms, respectively, and the values of  $\Delta$  are connected with the vibrational normal frequencies  $(\nu)$  through  $\Delta_b = (h/8\pi^2 v_b) \coth(hv_b/2kT)$ 

where h is Planck's constant, k Boltzmann's constant, and T the absolute temperature. Another set of relations, involving the force constants, is given below.

$$\varepsilon_1 = K_1 \Sigma_1 \tag{22}$$

$$\varepsilon_3 + \varepsilon_4 = \tilde{K}_2 \tilde{\Sigma}_3 + \Phi \Sigma_4 + 2\Gamma \Sigma_{34} \tag{23}$$

$$\begin{aligned}
\varepsilon_1 &= K_1 \Sigma_1 \\
\varepsilon_3 + \varepsilon_4 &= K_2 \Sigma_3 + \Phi \Sigma_4 + 2\Gamma \Sigma_{34} \\
\varepsilon_3 \varepsilon_4 &= (K_2 \Phi - \Gamma^2)(\Sigma_3 \Sigma_4 - \Sigma_{34}^2)
\end{aligned} \tag{22}$$

The force constants, viz.  $K_1$ ,  $K_2$ ,  $\Phi$  and  $\Gamma$ , are based on the chosen symmetry coordinates (Table 1), and given in terms of usual valence force coordinates in Ref.<sup>9</sup> The values of  $\varepsilon$  are given by

$$\varepsilon_k = \frac{1}{2} h \nu_k \, \coth \, \left( h \nu_k / 2kT \right) \tag{25}$$

If the coefficients of the transformation connecting the symmetry coordinates (S) with the normal coordinates (Q), viz.

$$S_i = \sum_k L_{ik} Q_k \tag{26}$$

are known, the symmetrized mean-square amplitude matrix elements are simply obtained by the following expressions.

$$\begin{array}{ll} \Sigma_{1} = L_{11}^{2} \Delta_{1} & (27) \\ \Sigma_{3} = L_{33}^{2} \Delta_{3} + L_{34}^{2} \Delta_{4} & (28) \\ \Sigma_{4} = L_{43}^{2} \Delta_{3} + L_{44}^{2} \Delta_{4} & (29) \\ \Sigma_{34} = L_{33} L_{43} \Delta_{3} + L_{34} L_{44} \Delta_{44} & (30) \end{array}$$

$$\Sigma_3 = L_{33}^2 \Delta_3 + L_{34}^2 \Delta_4 \tag{28}$$

$$\Sigma_{\mathbf{A}} = L_{\mathbf{A}3}^2 \Delta_3 + L_{\mathbf{A}4}^2 \Delta_{\mathbf{A}} \tag{29}$$

$$\Sigma_{34} = L_{33}L_{43}\Delta_3 + L_{34}L_{44}\Delta_4 \tag{30}$$

The most important ones of the mean-square amplitude quantities here concerned, are the two designated  $\sigma$  and  $\tau^*$ , and defined by eqns. (1) and (8), respectively. They are given in terms of the symmetrized mean-square amplitude matrix elements by the expressions

$$\sigma = \frac{1}{3}(\Sigma_1 + 2\Sigma_3) \tag{31}$$

$$\tau^* = \Sigma_1 + \frac{1}{2}(\Sigma_3 + \frac{1}{3}\Sigma_4) - 3^{-\frac{1}{2}}\Sigma_{34}$$
 (32)

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Table 2. Transformation coefficients  $L_{ik}$  in (atomic weight units)- $\frac{1}{2}$  for boron trihalides.

$^{11}\mathrm{BF_3}$	$Q_{1}$	$Q_3$	$Q_{4}$
$S_1$	0.2294	0.4346	0.001331
$S_1 \ S_3 \ S_4$	• • •	$0.4346 \\ 0.5413$	$0.001331 \\ 0.5230$
$^{11}\mathrm{BCl_3}$	$Q_1$	$Q_3$	$Q_4$
$S_1$	0.1679	0.4051	-0.01749
$S_1 \atop S_3 \atop S_4$	• • •	$0.4051 \\ 0.5981$	$-0.01749 \\ 0.3681$
<sup>11</sup> BBr <sub>3</sub>	$Q_{1}$	$Q_3$	$Q_4$
$S_1$	0.1119	0.3856	0.006125
$S_1 \ S_3 \ S_4$	•••	$0.3856 \\ 0.6074$	$0.006125 \\ 0.2779$
<sup>11</sup> BI <sub>3</sub>	$Q_1$	$Q_3$	$Q_{4}$
$S_1$	0.08876		
$S_1\ S_3\ S_4$	• • •	$0.3789 \\ 0.6333$	$-0.02293 \\ 0.1765$

These quantities represent the mean-square amplitudes of vibration for the bonded and non-bonded interatomic distances, respectively, and would also be obtainable from electron-diffraction. See, e.g. Refs.<sup>10–15</sup>

### NUMERICAL COMPUTATIONS

In the numerical computations reported in the following, the three vibrational frequencies for  $^{11}$ BY<sub>3</sub> (Y = F,Cl,Br,I) together with  $v_3$  for  $^{10}$ BY<sub>3</sub> were used to determine the vibrational constants. The experimental data for BF<sub>3</sub> and BCl<sub>3</sub> have been taken from Ref.<sup>4</sup>, those for BBr<sub>3</sub> and BI<sub>3</sub> from Ref.<sup>5</sup>, and are quoted in a previous paper <sup>7</sup>. This paper <sup>7</sup> also contains the calculated force constants.

The coefficients  $L_{ik}$  defined by eqn. (26), for the trihalides of the <sup>11</sup>B isotope, are given numerically in Table 2. The resulting mean-square amplitude quantities  $\Sigma_1$ ,  $\Sigma_3$ ,  $\Sigma_4$  and  $\Sigma_{34}$  at the absolute zero point and the temperature 298 °K, are given in Table 3.

Molecule	Symbol	Mean-square amplitude quantity $T=0$ 298°K	
<sup>11</sup> BF <sub>3</sub>	Σ-	0.000999	0.001027
	$\Sigma_{\mathbf{a}}$	0.002191	0.002195
	$\overline{\Sigma}_{\bullet}^{s}$	0.013000	0.015103
	$egin{array}{c} oldsymbol{\varSigma_1} \ oldsymbol{\varSigma_3} \ oldsymbol{\varSigma_4} \ oldsymbol{\varSigma_{34}} \end{array}$	0.002753	0.002763
<sup>11</sup> BCl <sub>3</sub>	Σ'	0.001010	0.001242
.DO13	<u>~</u> 1	0.002921	0.001242
	$\frac{2}{\Sigma}^3$	0.002321 $0.015723$	0.002333 $0.024279$
	$egin{array}{c} oldsymbol{\varSigma_1} \ oldsymbol{\varSigma_3} \ oldsymbol{\varSigma_4} \ oldsymbol{\varSigma_{34}} \end{array}$	0.003835	0.003521
$^{11}\mathrm{BBr_3}$	ς.	0.000759	0.001297
	Σ <sub>1</sub>	0.003063	0.001297
	\(\frac{7}{5}\)	0.016386	0.003130
	$egin{array}{c} \Sigma_1 \ \Sigma_3 \ \Sigma_4 \ \Sigma_{34} \end{array}$	0.005011	0.005572
$^{11}\mathrm{BI_3}$	Σ	0.000699	0.001631
1013	$\Sigma_{-}^{1}$	0.003527	0.001031 $0.004051$
	$\frac{23}{\Sigma}$	0.003327	0.032465
	$egin{array}{c} \Sigma_1 \ \Sigma_3 \ \Sigma_4 \ \Sigma_{34} \end{array}$	0.014000 $0.005064$	0.0032403

Table 3. Mean-square amplitude quantities in A<sup>2</sup> units for boron trihalides.

The mean amplitudes of vibration are in the present notation given by

$$u_{XY} = \sigma \frac{1}{2}, \qquad u_{YY} = (\tau^*)\frac{1}{2}$$
 (33)

and are obtainable according to eqns. (31) and (32). Their numerical values for the molecules here considered, have been given in Table 2 of Ref.<sup>7</sup>

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